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### Multi-scale modeling of turbulent oceanic flows

Vamsi K. Chalamalla UNC Chapel Hill

Edward Santilli Thomas Jefferson University, edward.santilli@jefferson.edu

Alberto Scotti UNC Chapel Hill

Masoud Jalali *UCSD* 

Sutanu Sarkar *UCSD* 

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Vamsi K Chalamalla UNC Chapel Hill

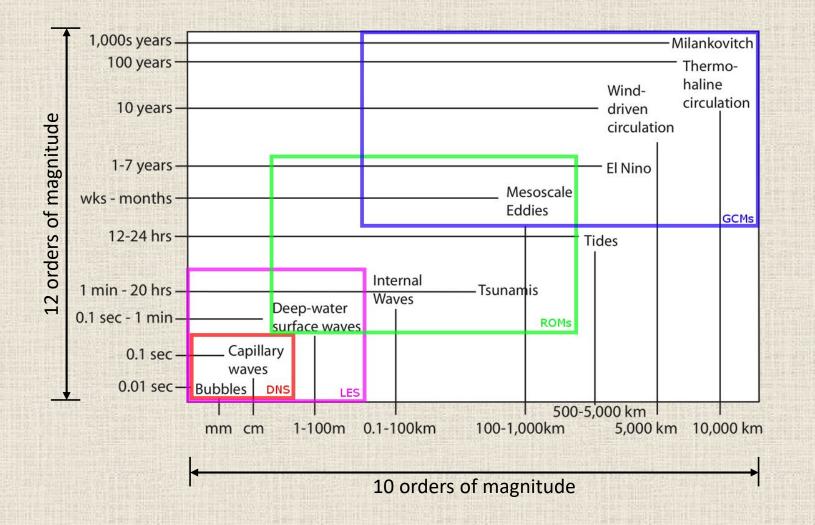
Edward Santilli Jefferson University

> Alberto Scotti UNC Chapel Hill

> > Masoud Jalali UCSD

Sutanu Sarkar UCSD Multi-scale modeling of turbulent oceanic flows

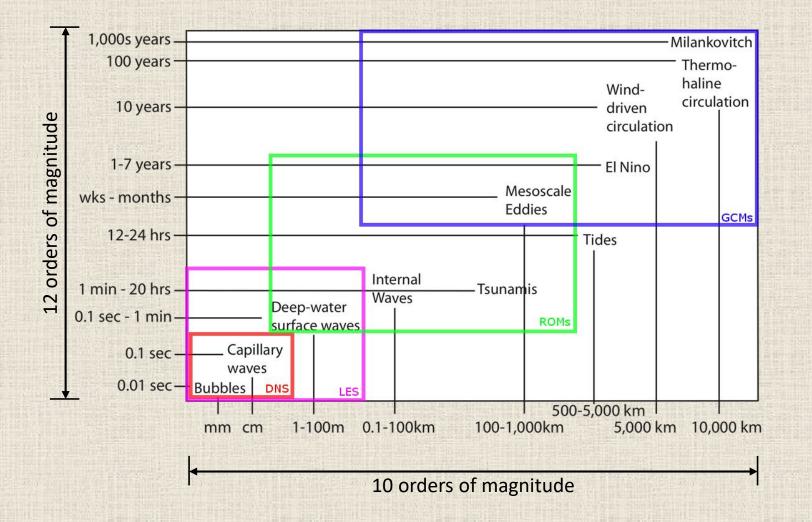
# **Modeling Challenges**



- LES Turbulence parameterized.
- DNS Turbulent scales resolved.
- Regional Ocean Models Captures nonlinear energy transfers. Requires a nonhydrostatic model.
- Global Circulation Models Small scale effects poorly parameterized.

Image: http://pordlabs.ucsd.edu/ltalley/sio210/introduction/index.html

# **Modeling Challenges**



LES – Turbulence parameterized.

DNS – Turbulent scales resolved.

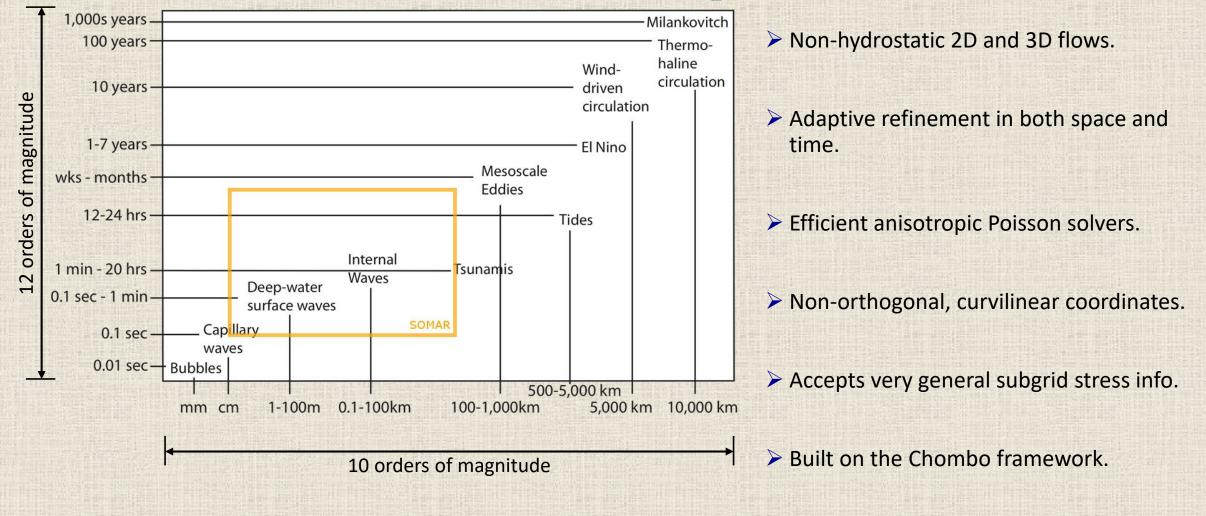
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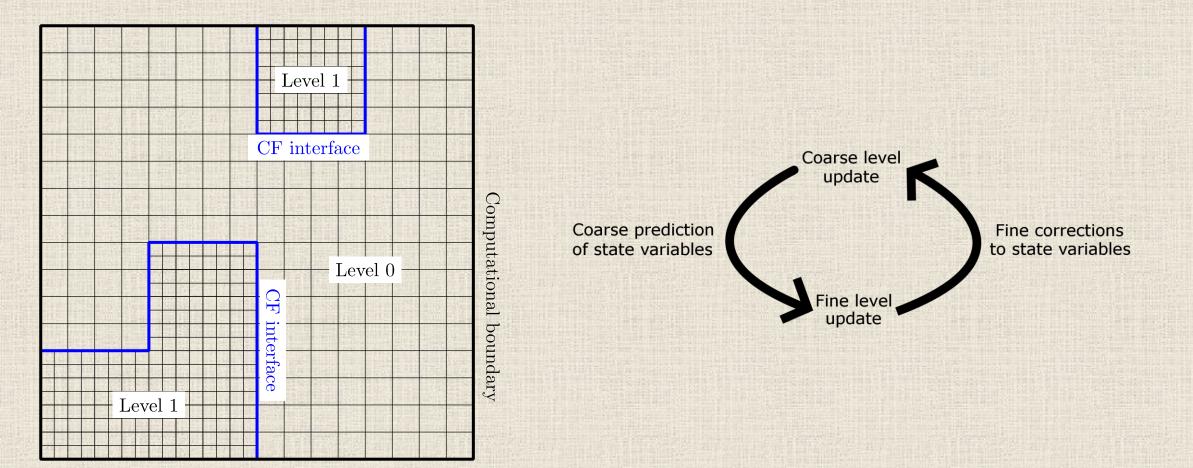
# SOMAR

### **Stratified Ocean Model with Adaptive Refinement**



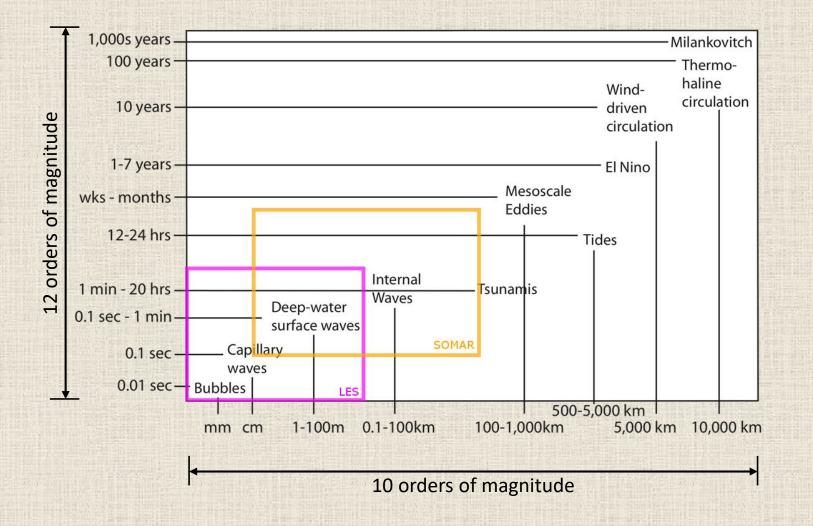
SOMAR: Santilli & Scotti (*Journal of Computational Physics*, 2011 & 2015) Chombo: https://commons.lbl.gov/display/chombo

## **SOMAR** Stratified Ocean Model with Adaptive Refinement



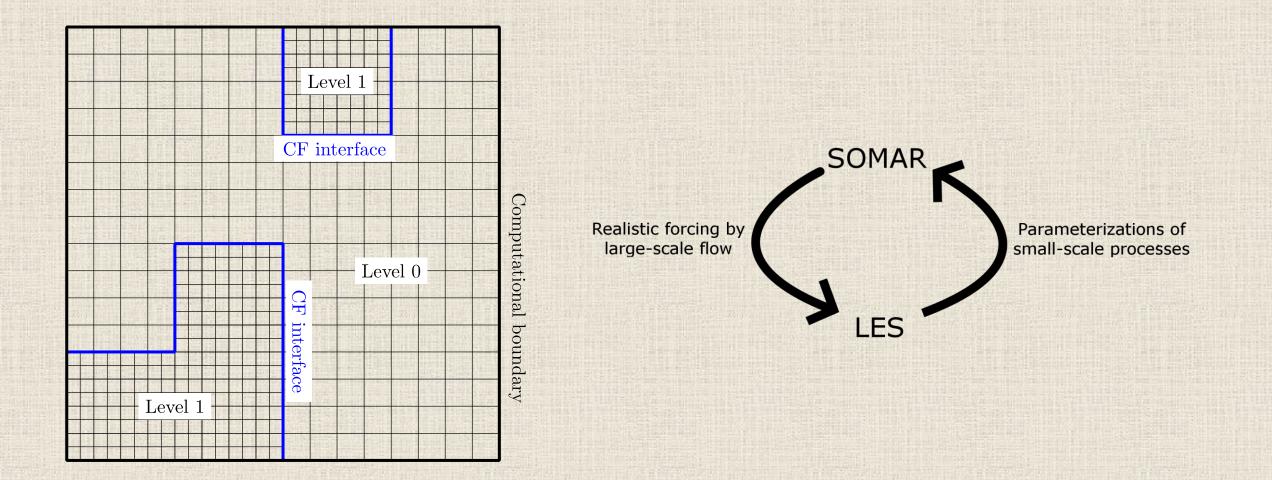
### Image: Chalamalla et al. (Ocean Modelling, 2017)

# **SOMAR-LES**



- > LES Turbulence parameterized.
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# **SOMAR-LES Algorithm**



# **SOMAR-LES Equations**

### **Coarse-grid equations**

 $\nabla \cdot \mathbf{u}$ 

### **Fine-grid equations**

$$\nabla \cdot \mathbf{u} = 0, \qquad \nabla \cdot \overline{\mathbf{u}} = 0$$

$$\frac{D\mathbf{u}}{Dt} = \nu \nabla^2 \mathbf{u} - \nabla \cdot \overline{\tau} - -f\hat{\mathbf{k}} \times \mathbf{u} + b^* \hat{\mathbf{k}} - \nabla p^* \qquad \qquad \frac{D\overline{\mathbf{u}}}{Dt} = \nu \nabla \nabla \frac{D\overline{\mathbf{u}}}{Dt} = \nu \nabla \frac{D\overline{\mathbf{u}}}{Dt} = \nu \nabla \frac{D\overline{\mathbf{u}}}{Dt} = \kappa \nabla \frac{D\overline{\mathbf{u}}}{Dt} = \kappa \nabla \frac{D\overline{\mathbf{u}}}{Dt} = \kappa \nabla \frac{D\overline{\mathbf{u}}}{Dt} = \kappa \nabla \frac{D\overline{\mathbf{u}}}{T} = \frac{D$$

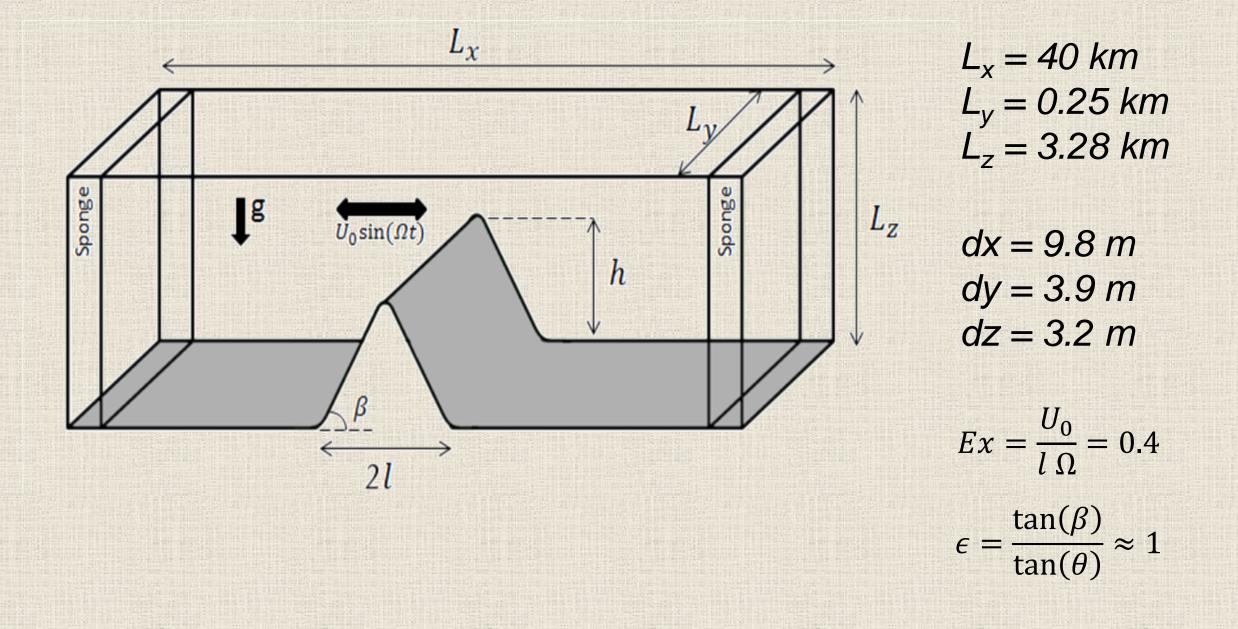
 $abla^2 \overline{\mathbf{u}} - 
abla \cdot oldsymbol{ au} - f \hat{\mathbf{k}} imes \overline{\mathbf{u}} + \overline{b}^* \hat{\mathbf{k}} - 
abla p^*$  $abla^2 \overline{b}^* - 
abla \cdot oldsymbol{\lambda} + \overline{w} N^2$  $\left.\begin{array}{c} -2\nu_{sgs}\overline{S}\\ -\kappa_{sgs}\nabla\overline{b}^{*}\end{array}\right\}$  $(\Delta)^{2}\sqrt{2\overline{S}:\overline{S}}$ 

$$\begin{aligned} \mathbf{ucros} : \nu_{sgs} &= 0.0014 \, C_K^{-3/2} \,\overline{\Delta} \, \sqrt{F_2} \\ \overline{F_2} &= \frac{1}{6} \left( \left| \left| \tilde{\bar{u}}_{i+1,j,k} - \tilde{\bar{u}}_{i,j,k} \right| \right|^2 + \left| \left| \tilde{\bar{u}}_{i-1,j,k} - \tilde{\bar{u}}_{i,j,k} \right| \right|^2 \right. \\ &+ \left| \left| \tilde{\bar{u}}_{i,j+1,k} - \tilde{\bar{u}}_{i,j,k} \right| \right|^2 + \left| \left| \tilde{\bar{u}}_{i,j-1,k} - \tilde{\bar{u}}_{i,j,k} \right| \right|^2 \\ &+ \left| \left| \tilde{\bar{u}}_{i,j,k+1} - \tilde{\bar{u}}_{i,j,k} \right| \right|^2 + \left| \left| \tilde{\bar{u}}_{i,j,k-1} - \tilde{\bar{u}}_{i,j,k} \right| \right|^2 \right) \end{aligned}$$

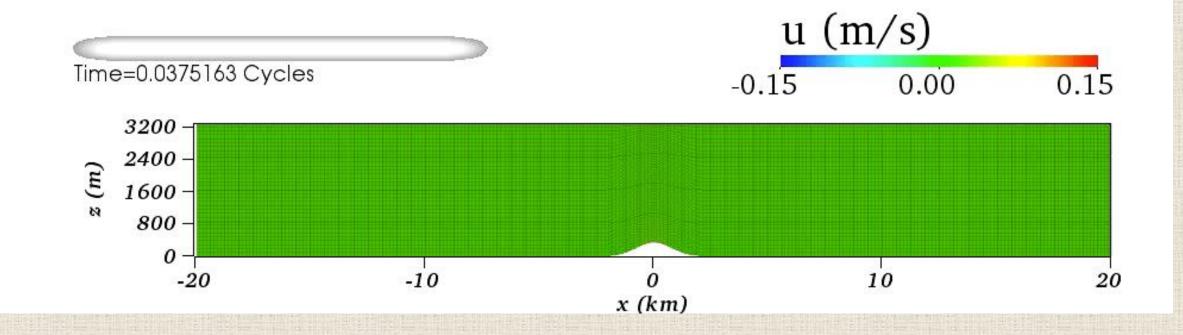
For more details: Chalamalla et al. (Ocean Modelling, 2017)

# **SOMAR-LES Simulations**

# **Internal tide generation with AMR**



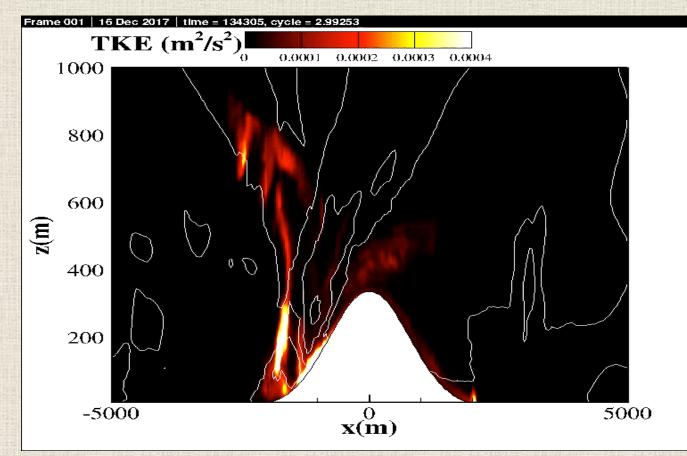
## **Internal tide generation with AMR**



Refinement only in the localized regions near the bathymetry.

 $\geq$  Refinement is done based on the gradient Richardson number Ri<sub>g</sub> < 0.25.

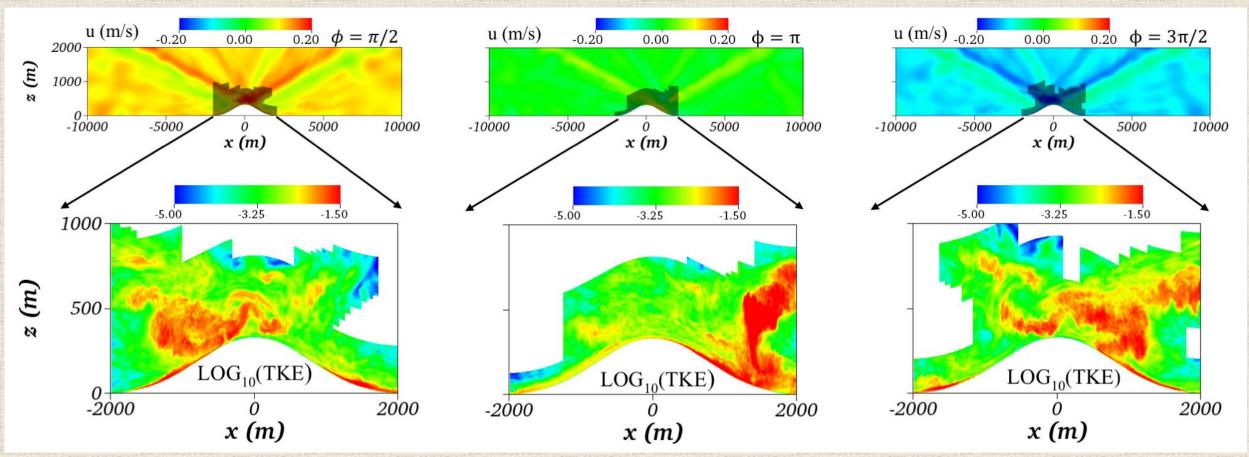
# Internal tide generation with AMR



> Turbulence is intermittent in both space and time.

Ideal problem for AMR, however the locations of turbulence needs to be predicted accurately.

# **Adaptive refinement + Subgrid scale model**



> Fine level grid adapts following the turbulent overturns.

> Turbulent overturns as tall as 500 m are found at certain phases.

> Need to a better job in predicting where the fine level grid is required.

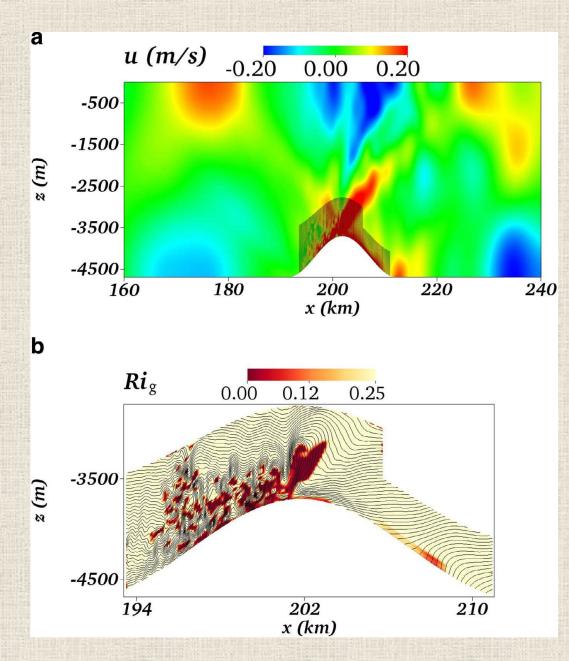
# **Model validation**

| Model     | С     | M     | q = 1-M/C | Р     |
|-----------|-------|-------|-----------|-------|
| SOMAR-LES | 0.736 | 0.637 | 0.13      | 0.036 |
| LES       | 0.721 | 0.612 | 0.15      | 0.031 |

Baroclinic energy budget and turbulent statistics compare well with previous numerical studies.

- Residual for baroclinic energy budget is less than 1%
- > Fine level grid occupies less than 2% of the total computational domain.
- > Total computational cost is just 10% of the single level grid solver.

## Low mode wave scattering

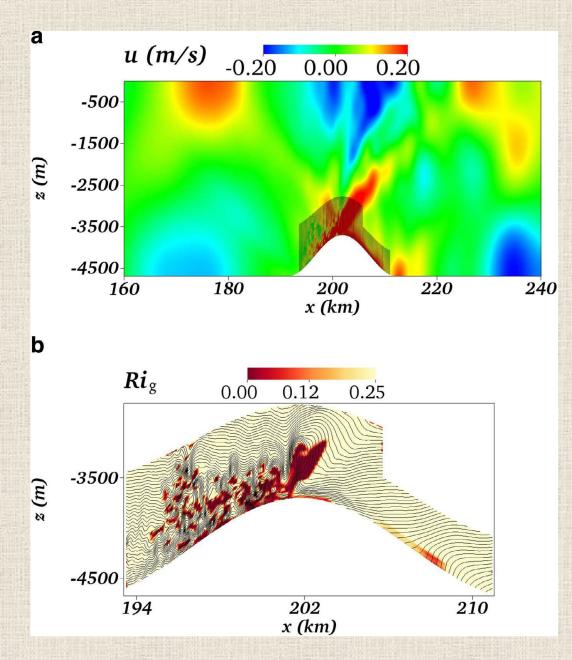


| $L_{x} = 446 \ km$ | dx = 109 m |
|--------------------|------------|
| $L_{y} = 127  km$  | dy = 125 m |
| $L_z = 4.7  km$    | dz = 18 m  |

*Topographic width* = 21 km

- Interaction of low-mode internal wave with an isolated bathymetry results in the generation of higher modes due to nonlinear interaction with the sloping bottom.
- Fine grids exist where Gradient Richardson number is < 0.25.</p>

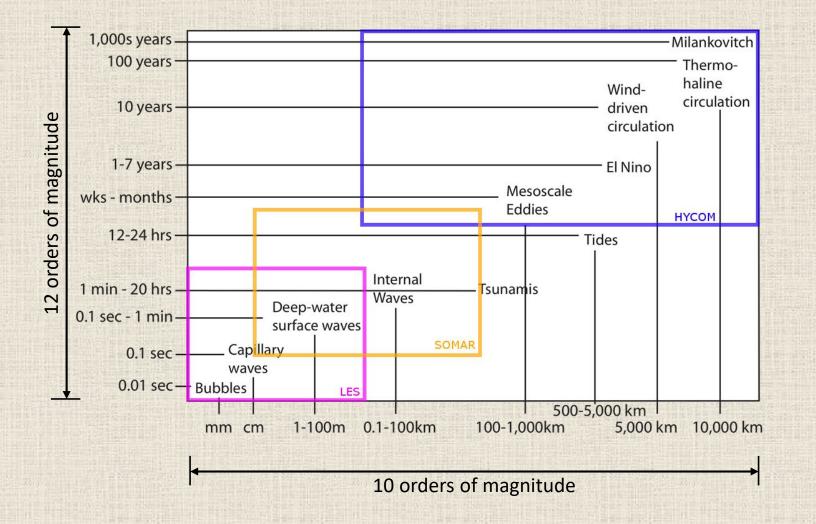
# Low mode wave scattering



### **SOMAR-LES vs. MITgcm**

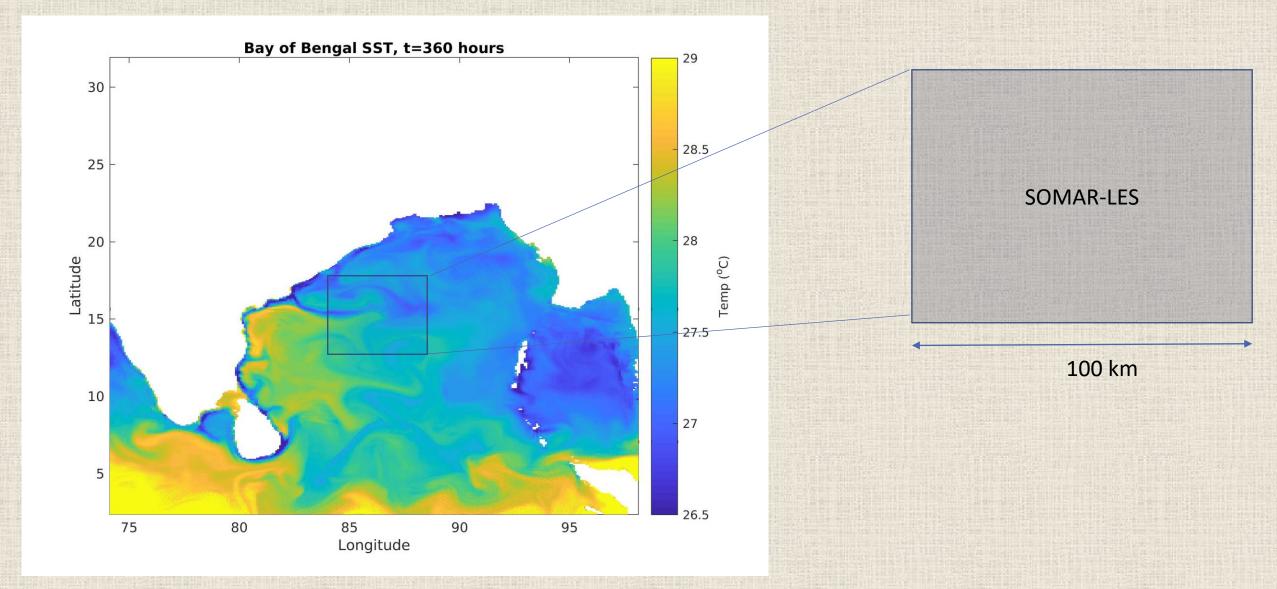
- SOMAR-LES used 3x finer resolution than MITgcm.
- > MITgcm used 4x more CPU-hours than SOMAR-LES.
- Reflected and transmitted energy in close agreement with linear theory.
- LES produces less dissipation than MITgcm's Thorpe-scale based eddy viscosity model.
- Work in progress: Add more levels of refinement + apply LES on more refined grid.

# **HYCOM-SOMAR-LES**



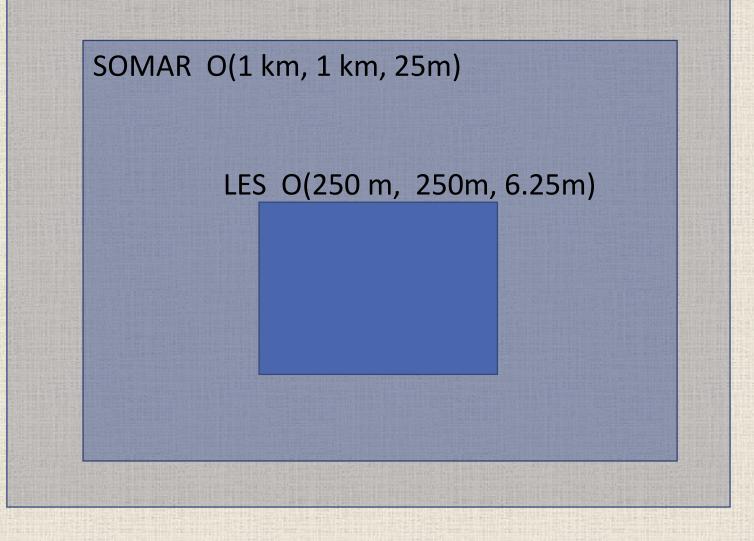
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**HYCOM-SOMAR-LES** 



### **Future work with HYCOM-SOMAR-LES**

### HYCOM O(4 km, 4 km, 100m)



To study ocean-atmosphere interactions in the Bay of Bengal region by coupling this with atmospheric model.

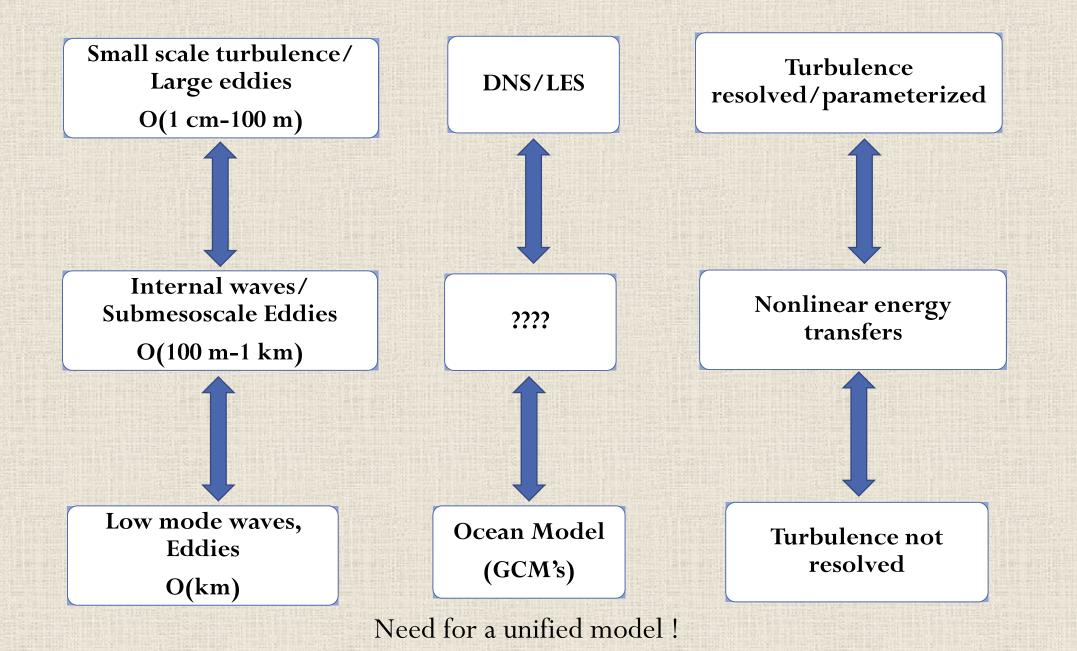
Achieve finer resolution, O(100m) in the horizontal and O(5m) in the vertical, to resolve the submesoscale eddies and their energy transfer to turbulence.

With finer resolution, we hope to numerically model features which are observed in the field experiments and absent in the global ocean simulations.

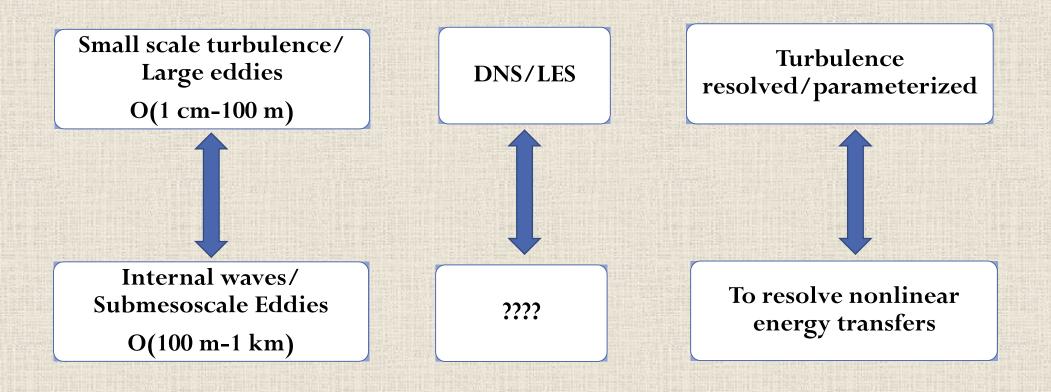
Address challenges that exist during exchange of information at interfacial boundaries.

# Thank you!

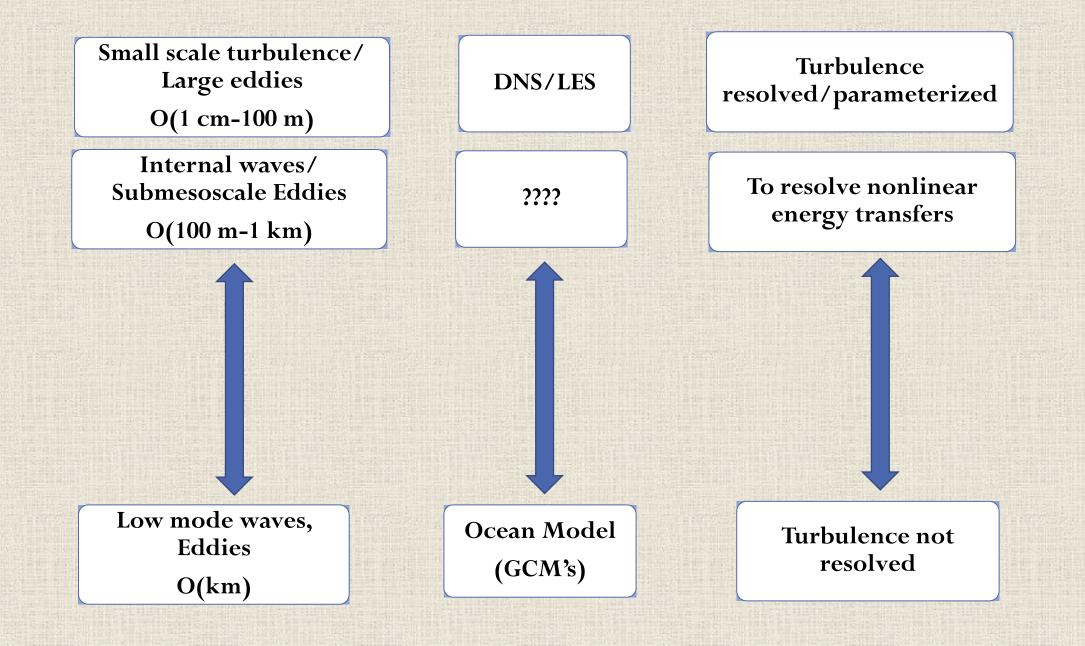
# Why multi-scale modeling?



# Why multi-scale modeling?



# Why multi-scale modeling?



# Large eddy simulations ..

Very expensive to resolve all the scales of the flow.
 Smallest length scales are removed via low pass filtering of Navier-Stokes equations.
 Effect of the unresolved scales are modeled via sub-grid scale modeling.

$$\nabla . \bar{u} = 0$$

$$\begin{split} \frac{\partial \overline{u}}{\partial t} &+ \overline{u}. \nabla \overline{u} = -\nabla p^* + \nu \nabla^2 \overline{u} - \frac{g}{\rho_0} \rho^* \hat{z} + \vec{F}_{ext} - \nabla. \tau \\ &\frac{\partial \rho^*}{\partial t} + \overline{u}. \nabla \rho^* = \kappa \nabla^2 \rho^* - \overline{w} \frac{d\rho^b}{dz} - \nabla. \lambda \end{split}$$

# **Eddy viscosity models**

Subgrid-scale stress tensor & density flux are defined as ...

$$\tau_{ij} = -2 \, \nu_{sgs} \, S_{ij}$$

$$\lambda_j = -\kappa_{sgs} \frac{\partial b^*}{\partial x_j}$$

Subgrid-scale viscosity & diffusivity is then computed using a variety of models.

➤Smagorinsky model

$$\nu_{sgs}(\boldsymbol{x},t) = (C_s \Delta^2) |S|$$

# **SOMAR-LES Equations**

### **Coarse-grid equations**

$$\begin{aligned} \nabla \cdot \mathbf{u} &= 0, \\ \frac{D\mathbf{u}}{Dt} &= \nu \nabla^2 \mathbf{u} - \nabla \cdot \boldsymbol{\tau} - -f\hat{\mathbf{k}} \times \mathbf{u} + b^* \hat{\mathbf{k}} - \nabla p^* \\ \frac{Db^*}{Dt} &= \kappa \nabla^2 b^* - \nabla \cdot \boldsymbol{\lambda} + wN^2 \end{aligned}$$

### **Fine-grid equations**

$$\nabla \cdot \overline{\mathbf{u}} = 0$$
  

$$\frac{D\overline{\mathbf{u}}}{Dt} = \nu \nabla^2 \overline{\mathbf{u}} - \nabla \cdot \boldsymbol{\tau} - f \hat{\mathbf{k}} \times \overline{\mathbf{u}} + \overline{b}^* \hat{\mathbf{k}} - \nabla p^*$$
  

$$\frac{D\overline{b}^*}{Dt} = \kappa \nabla^2 \overline{b}^* - \nabla \cdot \boldsymbol{\lambda} + \overline{w} N^2$$
  

$$\boldsymbol{\tau} = -2\nu_{sgs} \overline{S}$$
  

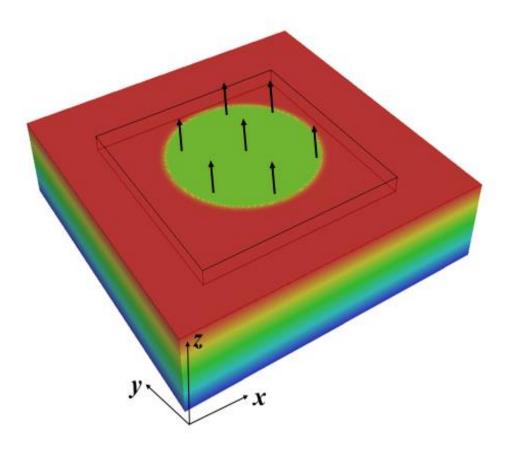
$$\boldsymbol{\lambda} = -\kappa_{sgs} \nabla \overline{b}^*$$
  

$$\nu_{sgs} = \left(C_S \sqrt[3]{\Delta x \, \Delta y \, \Delta z}\right)^2 \sqrt{2\overline{S} : \overline{S}}$$

Add Ducros!

Chalamalla et al. (Ocean Modelling, 2017)

# **Schematic**



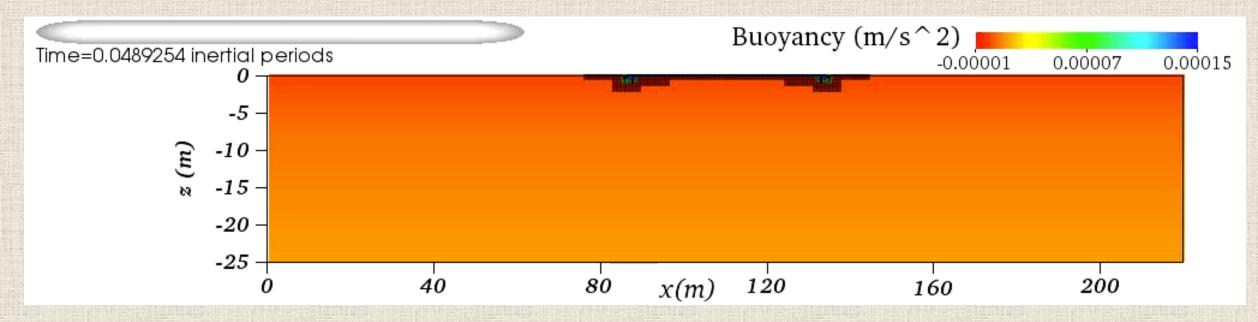
Surface cooling is applied by adding forcing term to r.h.s of the density equation.

Two-level grid with adaptive refinement is used to dynamically resolve convective plumes.

 $Rayleigh number = \frac{B_0 H^4}{\kappa^2 \nu}$ 

Rossby number =  $\frac{B_0^{1/2}}{f^{3/2}H}$ 

# **SOMAR-LES** Animations



Localized refinement only in the regions of turbulence based on gradient Richardson number criteria Ri<sub>g</sub> < 0</p>

>Add some points .. About efficiency of SOMAR in these kind of problems ..